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Weighting of single isomorphous replacement Fouriers. By R.H.Stanford, Jr., Gates and Crellin Laboratories of Chemistry,\* California Institute of Technology, Pasadena, California 91109, U.S.A.

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If an electron density map is calculated, using as the phase angle,  $\alpha_h$ , the average value of the pair of angles  $(\alpha_h + \delta_h)$  and  $(\alpha_h + \delta_h)$  derived from a single isomorphous replacement procedure, then an upper bound in the error of the resulting map, derived from the 'Cauchy inequality', will be minimized when the coefficient,  $|F_h|$ , in the Fourier summation is wieghted by the cosine of  $\delta_h$ .

The problems of using a single pair of isomorphous crystals in a structure determination have been discussed by Kartha (1961) and Blow & Rossmann (1961). The use of a single pair of isomorphous crystals results in an ambiguous pair of phase angles which may be represented as  $\alpha_h + \delta_h$  and  $\alpha_h - \delta_h$ ; only the sign is unknown, and  $\delta_h$  is restricted to the range 0 to  $\frac{1}{4}$ . Both papers cited above suggest the use of  $|F(h)| \cos 2\pi \delta_h$  as the coefficients and  $\alpha_h$  as the phase angles in the Fourier summation as a means of circumventing the problem of choosing the sign of  $\delta_h$ . This note lends theoretical support to this suggestion.

The true electron density is given by the expression:

$$\rho(\mathbf{r}) = \frac{1}{V} \sum_{\mathbf{h}} |F(\mathbf{h})| \exp \left[2\pi i (\alpha_{\mathbf{h}} \pm \delta_{\mathbf{h}})\right] \exp \left(-2\pi i \mathbf{h} \cdot \mathbf{r}\right), \quad (1)$$

where the correct sign for  $\delta_h$  should be chosen for each term of the summation, but is, by the nature of the procedure, not known. To stress this fact we carry the  $\pm$  indication until it disappears naturally near the final step of the analysis.

Assuming a phase angle equal to  $\alpha_h$  (average of the ambiguous pair), an approximate electron density may be defined by:

$$\rho'(\mathbf{r}) = \frac{1}{V} \sum_{\mathbf{h}} w_{\mathbf{h}} |F(\mathbf{h})| \exp(2\pi i \alpha_{\mathbf{h}}) \exp(-2\pi i \mathbf{h} \cdot \mathbf{r}), \quad (2)$$

where  $w_h$  is a weighting factor, assumed to be real, and is chosen in a manner described below to limit the difference between  $\rho(\mathbf{r})$  and  $\rho'(\mathbf{r})$ . Although it may not be possible, through the use of a simple weighting factor, to minimize the difference, it is possible to define a least upper bound.

Subtracting equation (2) from equation (1) yields:

$$\rho(\mathbf{r}) - \rho'(\mathbf{r}) = \frac{1}{V} \sum_{\mathbf{h}} |F(\mathbf{h})| \exp(2\pi i \alpha_{\mathbf{h}}) \exp(-2\pi i \mathbf{h} \cdot \mathbf{r})$$
$$[\exp(\pm 2\pi i \delta_{\mathbf{h}}) - w_{\mathbf{h}}]. \tag{3}$$

Cauchy's (1897) inequality may be cast in the following form:

$$\left|\sum_{\mathbf{r}} a_{\mathbf{r}} b_{\mathbf{r}}\right|^2 \leq \sum_{\mathbf{r}} |a_{\mathbf{r}}|^2 \cdot \sum_{\mathbf{r}} |b_{\mathbf{r}}|^2 \cdot$$

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Applying this inequality to equation (3) yields:

$$|\rho(\mathbf{r}) - \rho'(\mathbf{r})|^2 \le \frac{1}{V^2} \sum_{\mathbf{h}} ||F(\mathbf{h})| \exp(2\pi i \alpha_{\mathbf{h}}) \exp(-2\pi i \mathbf{h} \cdot \mathbf{r})|^2$$

$$\times \sum_{\mathbf{h}} |\exp(\pm 2\pi i \delta_{\mathbf{h}}) - w_{\mathbf{h}}|^2. \tag{4}$$

The first summation contains only known quantities and is, therefore, fixed. All terms,  $t_b$ , of the second summation are nonnegative real numbers:

$$t_{\mathbf{h}} = |\exp(\pm 2\pi i \delta_{\mathbf{h}}) - w_{\mathbf{h}}|^2$$
.

To establish a least upper bound for the squared error in the electron density,  $|\rho(\mathbf{r}) - \rho'(\mathbf{r})|^2$ , it is only necessary to minimize each term,  $t_h$ . Expanding  $t_h$  gives:

$$t_h = 1 + w_h^2 - 2w_h \cos 2\pi \delta_h$$
.

Note that this expansion involves the product of a number and its complex conjugate, leading to the disappearance of the sign ambiguity. To minimize  $t_h$ , we differentiate with respect to  $w_h$  and set the result equal to zero, yielding:

$$\frac{\partial t_{\mathbf{h}}}{\partial w_{\mathbf{h}}} = 2w_{\mathbf{h}} - 2\cos 2\pi \delta_{\mathbf{h}} = 0,$$

0

$$w_{\mathbf{h}} = \cos 2\pi \delta_{\mathbf{h}}$$
.

The use of this weighting factor, therefore, reduces the maximum error in the approximate electron density [equation (2)] to its least value. Blow & Crick (1959) arrived at a similar result by a somewhat different route.

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